

“Hidden” Momentum in an Isolated Brick?

Kirk T. McDonald

Joseph Henry Laboratories, Princeton University, Princeton, NJ 08544

(February 14, 2020)

1 Problem

The term “hidden” momentum was popularized by Shockley [1] in considerations of an electromechanical example, and essentially all subsequent use of this term has been for such examples, where one considers the system to consist of matter plus electromagnetic fields.

A definition of “hidden” momentum has been proposed by Daniel Vanzella [2] (see also [3]) which can be applied to mechanical systems as well, where a subsystem has a specified volume and can interact with the rest of the system via contact forces and/or transfer of mass/energy across its surface (which can be in motion),

$$\mathbf{P}_{\text{hidden}} \equiv \mathbf{P} - M\mathbf{v}_{\text{cm}} - \oint_{\text{boundary}} (\mathbf{x} - \mathbf{x}_{\text{cm}}) (\mathbf{p} - \rho\mathbf{v}_b) \cdot d\mathbf{Area} = - \int \frac{f^0}{c} (\mathbf{x} - \mathbf{x}_{\text{cm}}) d\text{Vol}, \quad (1)$$

where \mathbf{P} is the total momentum of the subsystem, $M = U/c^2$ is its total “mass”, U is its total energy, c is the speed of light in vacuum, \mathbf{x}_{cm} is its center of mass/energy, $\mathbf{v}_{\text{cm}} = d\mathbf{x}_{\text{cm}}/dt$, \mathbf{p} is its momentum density, $\rho = u/c^2$ is its “mass” density, u is its energy density, \mathbf{v}_b is the velocity (field) of its boundary, and

$$f^\mu = \frac{\partial T^{\mu\nu}}{\partial x^\nu}, \quad (2)$$

is the 4-force density due to the subsystem, with $T^{\mu\nu}$ being the stress-energy-momentum 4-tensor of the subsystem.

Consider an isolated (unstressed) brick in an inertial frame where it has constant velocity $v \hat{\mathbf{x}}$ parallel to, say, its longest dimension l . In this frame the brick has mass m . The total “hidden” momentum of this brick is zero, so consider a partition of the brick into two subsystems by an imaginary, moving surface, $x_{\text{boundary}} = a + ut$ for $0 < a < l$, while the brick extends over $vt < x < l + vt$. What is the “hidden” momentum in the portion of the brick at $0 < x < a$ at time $t = 0$?

2 Solution

Labeling the subsystem at $vt < x < a + ut$ by the subscript a , its time-dependent mass is,

$$m_a = \frac{a + (u - v)t}{l} m. \quad (3)$$

The time-dependent momentum of this subsystem is,

$$\mathbf{P}_a = m_a v \hat{\mathbf{x}}. \quad (4)$$

The center of mass of the subsystem has x -coordinate (for times when the boundary surface is within the brick),

$$x_{a,\text{cm}} = vt + \frac{a + (u - v)t}{2}, \quad (5)$$

and the velocity of the center of mass is,

$$\mathbf{v}_{a,\text{cm}} = \frac{\mathbf{u} + \mathbf{v}}{2}. \quad (6)$$

Thus,

$$\mathbf{P}_a - m_a \mathbf{v}_{a,\text{cm}} = m_a \left(\mathbf{v} - \frac{\mathbf{u} + \mathbf{v}}{2} \right) = \frac{m_a(\mathbf{v} - \mathbf{u})}{2}, \quad (7)$$

which is nonzero for unless $u = v$.

We now evaluate the “hidden” momentum at time $t = 0$ according to the first form of eq. (1). The boundaries in x of subsystem a at this time are at $x = 0$ and a , with velocities v and u , while $x_{a,\text{cm}}(0) = a/2$.

$$\begin{aligned} \mathbf{P}_{a,\text{hidden}} &= \mathbf{P}_a - m_a \mathbf{v}_{a,\text{cm}} + \left(0 - \frac{a}{2}\right) \frac{m}{l} (\mathbf{v} - \mathbf{v}) - \left(a - \frac{a}{2}\right) \frac{m}{l} (\mathbf{v} - \mathbf{u}) \\ &= \frac{ma}{2l} (\mathbf{v} - \mathbf{u}) - \frac{ma}{2l} (\mathbf{v} - \mathbf{u}) = 0. \end{aligned} \quad (8)$$

To use the second form of eq. (1), we note that $T^{00} = mc^2/Al$ and $T^{0x} = mcv/Al$ are constant within subsystem a , whose cross-sectional area is A , while $T^{0y} = 0 = T^{0z}$ everywhere. The time component f^0 of the 4-force density (2) is,

$$f^0 = \partial_0 T^{00} + \partial_i T^{0i} = \frac{\partial T^{00}}{\partial ct} + \frac{\partial T^{0x}}{\partial x}. \quad (9)$$

The time dependence of T^{00} of subsystem a is due only to the presence of its moving boundaries, $x = vt$ and $x = a + ut$. Near the “left” boundary, $T^{00}(x, t) = T^{00}(x - vt)$, such that $\partial T^{00}/\partial ct = -(v/c)\partial T^{00}/\partial x$, while near the “right” boundary, $T^{00}(x, t) = T^{00}(x - a - ut)$, such that $\partial T^{00}/\partial ct = -(u/c)\partial T^{00}/\partial x$.

It suffices to complete the calculation of $\mathbf{P}_{a,\text{hidden}}$ for $t = 0$, as the result should be independent of time. At $t = 0$, f^0 is nonzero only at/near the boundaries $x = 0$ and $x = a$, so we can split the integration to that over $0 < x < b < a$ and $b < x < a$,

$$\begin{aligned} P_{\text{hidden},x}(t = 0) &= - \int \frac{f^0}{c} (x - x_{\text{cm}}) d\text{Vol} = - \frac{A}{c} \int_0^a dx f^0 \left(x - \frac{a}{2}\right) \\ &= \frac{Av}{c^2} \int_0^b dx \frac{\partial T^{00}}{\partial x} \left(x - \frac{a}{2}\right) + \frac{Au}{c^2} \int_b^a dx \frac{\partial T^{00}}{\partial x} \left(x - \frac{a}{2}\right) - \frac{A}{c} \int_0^a dx \frac{\partial T^{0x}}{\partial x} \left(x - \frac{a}{2}\right) \\ &= \frac{Av}{c^2} \left[xT^{00} \Big|_0^b - \int_0^b dx T^{00} - \frac{a}{2} [T^{00}(b) - T^{00}(0)] \right] \\ &\quad + \frac{Au}{c^2} \left[xT^{00} \Big|_b^a - \int_b^a dx T^{00} - \frac{a}{2} [T^{00}(a) - T^{00}(b)] \right] \\ &\quad - \frac{A}{c} \left[xT^{0x} \Big|_0^a - \int_0^a dx T^{0x} - \frac{a}{2} [T^{0x}(a) - T^{0x}(0)] \right] = 0, \end{aligned} \quad (10)$$

in agreement with eq. (8), taking T^{00} and T^{0x} to have values their nonzero, constant values within the interval $0 \leq x \leq a$. and zero outside this.¹

Thus, according to the calculations (9) and (10), there is no “hidden” momentum in the “all-mechanical” example of an isolated brick, or in subsystems of it defined by moving partitions.

As noted in sec. VI of [4], “hidden” momentum is associated with (sub)systems that have internal motion when “at rest”, which is not the case for an isolated brick.

References

- [1] W. Shockley and R.P. James, “Try Simplest Cases” Discovery of “Hidden Momentum” Forces on “Magnetic Currents”, Phys. Rev. Lett. **18**, 876 (1967),
http://physics.princeton.edu/~mcdonald/examples/EM/shockley_prl_18_876_67.pdf
- [2] D. Vanzella, Private communication, (June 29, 2012).
- [3] K.T. McDonald, *On the Definition of “Hidden” Momentum* (July 9, 2012),
<http://physics.princeton.edu/~mcdonald/examples/hiddendef.pdf>
- [4] D. Babson *et al.*, *Hidden momentum, field momentum, and electromagnetic impulse*, Am. J. Phys. **77**, 826 (2009),
http://physics.princeton.edu/~mcdonald/examples/EM/babson_ajp_77_826_09.pdf

¹An analysis of eq. (10) which invokes Heaviside step functions Θ , and Dirac delta functions δ , notes that in the frame where the rod has velocity v , the nonzero components of $T^{0\mu}$ can be written as,

$$T^{00} = \frac{mc^2}{Al} [\Theta(x - vt) - \Theta(x - a - ut)], \quad (11)$$

$$T^{0x} = \frac{mcv}{Al} [\Theta(x - vt) - \Theta(x - a - ut)], \quad (12)$$

where $\Theta(x) = 1$ for $x > 0$ and $= 0$ for $x < 0$. Then,

$$\frac{\partial T^{00}}{\partial ct} = -\frac{mc^2}{Al} \left[\frac{v}{c} \delta(x - vt) - \frac{u}{c} \delta(x - a - ut) \right], \quad (13)$$

$$\frac{\partial T^{0x}}{\partial x} = \frac{mcv}{Al} [\delta(x - vt) - \delta(x - a - ut)], \quad (14)$$

$$\begin{aligned} P_{\text{hidden},x}(t=0) &= - \int \frac{f^0}{c} (x - x_{\text{cm}}) d\text{Vol} = -\frac{A}{c} \int_0^a dx \left(\frac{\partial T^{00}}{\partial ct} + \frac{\partial T^{0x}}{\partial x} \right) (x - x_{\text{cm}}) \\ &= \frac{A}{c} \int_0^a dx \frac{mc^2}{Al} \left[\frac{v}{c} \delta(x) - \frac{u}{c} \delta(x - a) \right] (x - a/2) - \frac{A}{c} \int_0^a dx \frac{mcv}{Al} [\delta(x) - \delta(x - a)] (x - a/2) \\ &= \frac{ma}{l} \left(-u - \frac{v}{2} + \frac{u}{2} \right) + \frac{mav}{l} = \frac{ma(v - u)}{2l}. \end{aligned} \quad (15)$$

Note that the result of eq. (15) is the same as $\mathbf{P}_a - m_a \mathbf{v}_{a,\text{cm}}$ of eq. (7). Hence, if the boundary integral in the first form of eq. (1) were ignored, the two forms of that expression, according to calculations using delta functions in f^0 , would both lead to the same, nonzero “hidden” momentum in present example.

This author finds the delta functions in the expressions (13)-(14) for the 4-force density f^μ very unappealing physically, and so prefers the analysis in the main text that avoids them, with the implication that there is zero “hidden” momentum in the present example.